

Lecture 30: STABILITY OF TIME ITERATION

30

We seek to understand the apparent time step restriction for the Forward Euler method for the heat equation.

First consider a generic scalar differential equation

$$y'(t) = \lambda y(t)$$

for which the forward Euler method takes the form

$$y_{k+1} = y_k + \Delta t \lambda y_k = (1 + \Delta t \lambda) y_k$$

so that
$$\boxed{y_k = (1 + \Delta t \lambda)^k y_0}$$

Now the exact solution of the differential equation is $y(t) = e^{\lambda t} y(0)$, which conveys, $y(t) \rightarrow 0$, provided $\operatorname{Re}(\lambda) < 0$, i.e., if $\lambda = \alpha + i\beta$, then

$$\begin{aligned} e^{\lambda t} &= e^{(\alpha+i\beta)t} = e^{\alpha t} e^{i\beta t} \\ &= e^{\alpha t} (\cos \beta t + i \sin \beta t) \end{aligned}$$

$$\begin{aligned} \text{so } |e^{\lambda t}| &= |e^{\alpha t} (\cos \beta t + i \sin \beta t)| \\ &= e^{\alpha t} |\cos \beta t + i \sin \beta t| \\ &= e^{\alpha t} \underbrace{\sqrt{\cos^2 \beta t + \sin^2 \beta t}}_{=1} = e^{\alpha t}. \end{aligned}$$

So $e^{\lambda t} \rightarrow 0$ if $\operatorname{Re}(\lambda) = \alpha < 0$.

Does forward Euler mimic this behavior? 30.

$$y_k = (1 + \Delta t \lambda)^k y_0$$

Implies that $y_k \rightarrow 0$ provided $|1 + \Delta t \lambda| < 1$.

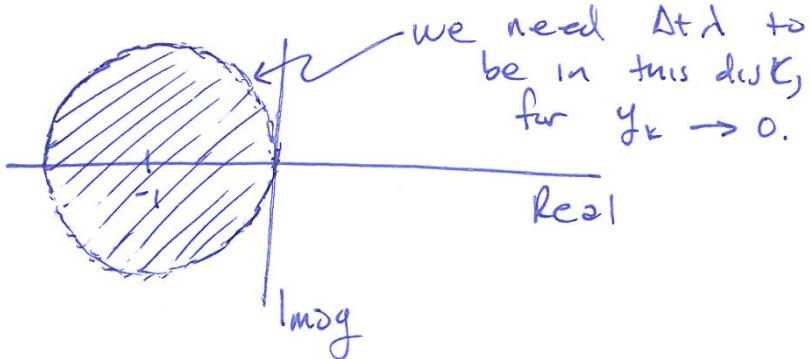
Note that this is a rather different condition than $\operatorname{Re}(\lambda) < 0$!

Consider real values of λ , as we have for eigenvalues of the heat discretization. If λ is negative but large in magnitude, we will need Δt to be very small for $|1 + \Delta t \lambda| < 1$.

Here is a helpful interpretation of this condition:

$\{z \in \mathbb{C} : |c - z| < 1\}$ = set of all complex numbers z for which $|c - z| < 1$
= set of all complex numbers z whose distance from c is less than one.

Thus $|1 + \Delta t \lambda| = |-1 - \Delta t \lambda| < 1$ is satisfied if $\Delta t \lambda$ is in the disc of radius 1 in the complex plane, centered at -1.



Now let us interpret this observation in light of the eigenvalue calculations in the last lecture. 30.3

The heat equation / Finite Element Method give

$$\dot{z}(t) = -M^{-1}K z(t)$$

So Forward Euler takes the form

$$z_{j+1} = z_j - \Delta t M^{-1} K z_j \Rightarrow z_j = (I - \Delta t M^{-1} K)^j z_0$$

where $z_j \approx z(j \Delta t)$.

In the last lecture we wrote

$$M^{-1}K = V \Lambda V^T = \sum_{k=1}^n \lambda_k v_k v_k^T$$

$$\text{where } \lambda_k = \frac{6}{h^2} \left(\frac{2 - \gamma_k}{4 + \gamma_k} \right), \quad \gamma_k = 2 \cos(k\pi/(N+1)), \\ h = \frac{1}{N+1}.$$

Forward Euler then becomes

$$\begin{aligned} z_j &= (I - \Delta t M^{-1} K)^j z_0 \\ &= (I - \Delta t V \Lambda V^T)^j z_0 \\ &= (V V^T - \Delta t V \Lambda V^T)^j z_0 \\ &= V (I - \Delta t \Lambda)^j V^T z_0 \\ &= V \begin{pmatrix} (1 - \Delta t \lambda_1)^j & & & \\ & \ddots & & 0 \\ & & \ddots & \\ & & & (1 - \Delta t \lambda_N)^j \end{pmatrix} V^T z_0 \\ &= \sum_{k=1}^N (1 - \Delta t \lambda_k)^j v_k v_k^T z_0 \end{aligned}$$

* The component of this sum that grows fastest is the one for which $|1 - \Delta t \lambda_k|$ is largest. This corresponds to λ_N when Δt violates the stability condition \Rightarrow the unstable solution will be dominated by v_N , the "smallest" eigenvector $vvvvv$.

Note:
 $VV^T = I$
 $V^T V = I$

So if we want $\lambda_j \rightarrow 0$ as $j \rightarrow \infty$, 30.

We need $|1 - \Delta t + \lambda_k| < 1$ for all $k=1, \dots, N$.

Since $\lambda_k \in (0, \frac{12}{h^2})$ for all $k=1, \dots, N$,

we should pick Δt sufficiently small that

$$\left| 1 - \underbrace{\Delta t \frac{12}{h^2}}_{>0} \right| < 1$$

Since $\Delta t \frac{12}{h^2} > 0$, for any Δt , $|1 - \Delta t \frac{12}{h^2}| < 1$,

Hence, Δt will be determined by the condition

$$\begin{aligned} -1 &< 1 - \frac{\Delta t \cdot 12}{h^2} \\ \Rightarrow \Delta t \frac{12}{h^2} &< 2 \Rightarrow \boxed{\Delta t < \frac{h^2}{6}} \end{aligned}$$

Notice the profound implication of this calculation:

If you want to double N (to get better accuracy in space), you reduce h by $\frac{1}{2}$, and you must then reduce Δt by $\underline{\underline{\frac{1}{4}}}$.

This severe stability constraint is called the CFL condition after its discoverers: Richard Courant, Kurt Friedrichs, and Hans Lewy.

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EACH DIFFERENT TIME INTEGRATOR LEADS TO
A DIFFERENT STABILITY CONSTRAINT.

AS AN EXAMPLE, CONSIDER THE BACKWARD EULER
METHOD. FOR THE GENERIC EQUATION $y'(t) = \lambda y(t)$,
NOTE THAT

$$y'(t) = \lim_{\Delta t \rightarrow 0} \frac{y(t+\Delta t) - y(t)}{\Delta t}$$

ALSO SUGGESTS THAT

$$0 = \lim_{\Delta t \rightarrow 0} y'(t+\Delta t) = \left(\frac{y(t+\Delta t) - y(t)}{\Delta t} \right)$$

HENCE, FOR FINITE Δt , APPROXIMATE

$$y(t+\Delta t) \approx y(t) + \Delta t y'(t+\Delta t)$$

WRITING $y_j \approx y(j \Delta t)$, THIS GIVES

$$y_{j+1} = y_j + \Delta t \lambda y_{j+1} \quad (y' = \lambda y)$$

$$\Rightarrow (1 - \Delta t \lambda) y_{j+1} = y_j$$

$$\Rightarrow y_{j+1} = \frac{1}{1 - \Delta t \lambda} y_j$$

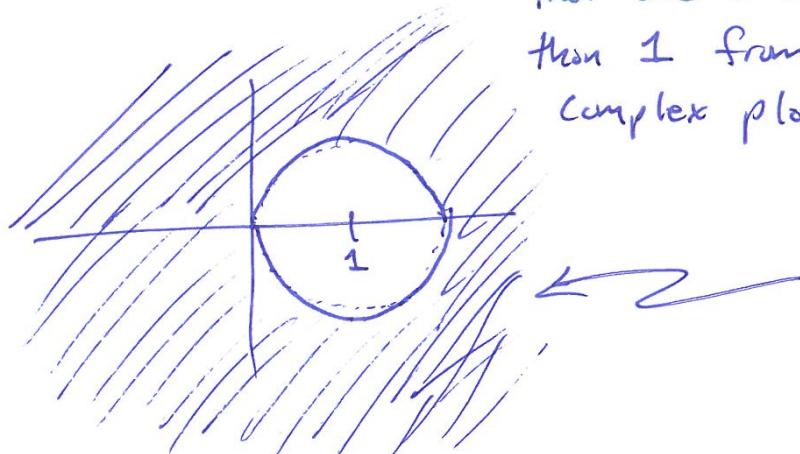
AND SO WE CAN ACCUMULATE j STEPS AS

$$y_j = \left(\frac{1}{1 - \Delta t \lambda} \right)^j y_0$$

NOW IF $\text{Re}(\lambda) < 0$, $y(t) \rightarrow 0$. IS THE SAME TRUE
FOR BACKWARD EULER?

$y_j \rightarrow 0$ if $\left| \frac{1}{1-\Delta t + \lambda} \right| < 1$, which implies 30.1

$|1 - \Delta t + \lambda| =$ set of all $\Delta t + \lambda$
that are a distance greater
than 1 from 1 in the
complex plane



So $y_j \rightarrow 0$
provided
 $\Delta t + \lambda$ is in
the shaded
region....

If $\operatorname{Re}(\lambda) < 0$, then for any Δt ,

$\operatorname{Re}(\Delta t + \lambda) < 0$, and so $\Delta t + \lambda$ is in the
shaded region.

⇒ BACKWARD EULER IS STABLE FOR ANY Δt .

(You will still adjust Δt to give an
accurate solution, i.e., to make the
approximation

$$y' \approx \frac{y(t+\Delta t) - y(t)}{\Delta t} \text{ accurate.}$$

For the heat equation, $\mathbf{z}'(t) = -M^{-1}K\mathbf{z}(t)$,

Backward Euler gives

$$\mathbf{z}_j = ((I + \Delta t M^{-1} K)^{-1})^j \mathbf{z}_0 = \sum_{k=1}^N \underbrace{\left(\frac{1}{1 + \Delta t \lambda_k} \right)^j}_{\lambda_k \in (0, \frac{12}{n^2})} V_k V_k^T \mathbf{z}_0$$

Since $\lambda_k > 0$, $\left| \frac{1}{1 + \Delta t \lambda_k} \right| < 1$ for all Δt .

Hence Backward Euler is stable for the heat equation for any $\Delta t > 0$. 30.7

CODA: A WORD ABOUT TIME-STEPPING IMPLEMENTATIONS
(See Lecture 35 for related detail.)

Forward Euler:

$$\bar{z}_{j+1} = \bar{z}_j - \Delta t M^{-1} K \bar{z}_j$$

Backward Euler:

$$\begin{aligned}\bar{z}_{j+1} &= (I + \Delta t M^{-1} K)^{-1} \bar{z}_j \\ &= (M^{-1} M + \Delta t M^{-1} K)^{-1} \bar{z}_j \\ &= (M^{-1} (M + \Delta t K))^{-1} \bar{z}_j \\ &= (M + \Delta t K)^{-1} M^{-1} \bar{z}_j \\ &= (M + \Delta t K)^{-1} M \bar{z}_j\end{aligned}$$

Both equations involve the inverse of one matrix, which can be implemented via LU factorization \Rightarrow similar work per step!

For more complicated equations (especially nonlinear equations) Backward Euler is rather more expensive per step than Forward Euler.