Convergence Theory for Iterative Eigensolvers

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setting for the talk

- ▶ Let $A \in \mathbb{C}^{n \times n}$ be a large square matrix, potentially non-Hermitian ($A \neq A^*$).
- Computing all eigenvalues of A is too expensive (and usually not needed).
- Thus we seek m le n distinguished eigenvalues relevant to our application (*largest, smallest, rightmost, etc.*)

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Projection Methods

 $\mathcal{V} \subset \mathbb{C}^n = k$ -dimensional subspace of \mathbb{C}^n , the projection subspace for **A** The columns of $\mathbf{V} \in \mathbb{C}^{n \times k}$ for an orthonormal basis for \mathcal{V} :

$$\mathbf{V}^*\mathbf{V} = \mathbf{I}, \qquad \mathbf{V}^*\mathbf{A}\mathbf{V} \in \mathbb{C}^{k \times k}$$

We hope *some* eigenvalues of V^*AV $\sigma(V^*AV) = \{\theta_1, \dots, \theta_k\}$ approximate *some* eigenvalues of **A**. $\sigma(\mathbf{A}) = \{\lambda_1, \dots, \lambda_n\}$ For example, $\theta_1 \approx \lambda_1, \dots, \theta_m \approx \lambda_m$ for some $1 \le m \le k$.

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This talk mainly describes established results for the deterministic case, with some thoughts from a RandNLA perspective along the way.

why consider $A \neq A^*$?

While Hermitian problems are common (SVD, quantum mechanics, etc.), many important applications lead to non-Hermitian problems – and subtler issues of spectral perturbation theory. Many examples: [Trefethen, E. 2005].

- atmospheric science
- fluid flow stability
- damped mechanical systems
- control theory
- data-driven modeling
- lasers
- ecology
- Markov chains





An Overview of

Projection-Based

Eigensolvers

projection-based eigensolvers

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Power method (minimal storage, easy to implement, can be slow) $\mathcal{V} = \text{span}\{\mathbf{A}^{p}\mathbf{x}\}$

$$\begin{split} \textbf{Subspace iteration} & (\text{more storage, subtler to implement, computes repeated eigs}) \\ \mathcal{V} = \text{Range}(\textbf{A}^p\textbf{X}) \text{ for } \textbf{X} \in \mathbb{C}^{n \times k} \end{split}$$

[Halko, Martinsson, Tropp 2011] et al.

 $\label{eq:variable} \begin{array}{ll} \mbox{Power method} & (\mbox{minimal storage, easy to implement, can be slow}) \\ \\ \mathcal{V} = \mbox{span}\{ \mathbf{A}^p \mathbf{x} \} \end{array}$

Subspace iteration (more storage, subtler to implement, computes multiple eigs) $\mathcal{V} = \text{Range}(\mathbf{A}^{p}\mathbf{X}) \text{ for } \mathbf{X} \in \mathbb{C}^{n \times b}$

 $\label{eq:relation} \begin{array}{ll} \mbox{Krylov subspace methods} & (\mbox{growing subspace dimension; higher powers of A}) \\ \\ \mathcal{V} = \mbox{span}\{x, Ax, A^2x, \dots, A^{k-1}x\} \end{array}$

Block Krylov methods(subspace dimension grows quickly: dim(\mathcal{V}) $\leq kb$) $\mathcal{V} = \mathsf{Range}([X \ AX \ A^2X \ \cdots \ A^{k-1}X])$ for $X \in \mathbb{C}^{n \times b}$ SVD: [Musco & Musco 2015], [Drineas et al. 2018]

Must balance benefit of large k with block size b, storage.

projection-based eigensolvers: krylov methods (extensions)

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Restarted Krylov (used in eigs: filter ψ improves starting vector) $\mathcal{V} = \operatorname{span}\{\psi(\mathbf{A})\mathbf{x}, \mathbf{A}\psi(\mathbf{A})\mathbf{x}, \mathbf{A}^{2}\psi(\mathbf{A})\mathbf{x}, \dots, \mathbf{A}^{k-1}\psi(\mathbf{A})\mathbf{x}\}$

$$\label{eq:very high degree polys, care needed} \begin{split} \textbf{Polynomial Preconditioned Krylov} & (very high degree polys, care needed) \\ \mathcal{V} = \text{span}\{\textbf{x}, \pi(\textbf{A})\textbf{x}, \pi(\textbf{A})^2\textbf{x}, \dots, \pi(\textbf{A})^{k-1}\textbf{x}\} \end{split}$$

$$\begin{split} \text{Shift-Invert Krylov} & (\text{used in eigs: ideal for eigenvalues near } \mu) \\ \mathcal{V} = \text{span}\{\mathbf{x}, (\mathbf{A} - \mu \mathbf{I})^{-1})\mathbf{x}, (\mathbf{A} - \mu \mathbf{I})^{-2}\mathbf{x}, \dots, (\mathbf{A} - \mu \mathbf{I})^{-(k-1)}\mathbf{x}\} \end{split}$$

Rational Krylov (helps for finding eigenvalues in a region) $\mathcal{V} = \operatorname{span}\{\mathbf{x}, (\mathbf{A} - \mu_1)^{-1}\mathbf{x}, (\mathbf{A} - \mu_2)^{-1}\mathbf{x}, \dots, (\mathbf{A} - \mu_{k-1}\mathbf{I})^{-1}\mathbf{x}\}$

preliminaries: spectral structure of A

• Distinct eigenvalues of A: $\lambda_1, \lambda_2, \ldots, \lambda_{\hat{n}}$

► Spectral projectors **P**_j and invariant subspaces U_j:

$$\mathbf{P}_j := rac{1}{2\pi i} \int_{\Gamma_j} (z\mathbf{I} - \mathbf{A})^{-1} \, \mathrm{d}z, \qquad \mathfrak{U}_j := \operatorname{Range}(\mathbf{P}_j),$$

 Γ_j is a contour in $\mathbb C$ containing λ_j but no other distinct eigenvalues.

• If $\mathbf{A} = \mathbf{A}^*$ and λ_j is simple with unit eigenvector \mathbf{u}_j , then $\mathbf{P}_j = \mathbf{u}_j \mathbf{u}_j^*$.

P_j is a projector onto the invariant subspace U_j, but P_j need not be an orthogonal projector when A ≠ A^{*}.

• The spectral projectors give a *resolution of the identity*: $\sum_{j=1}^{\hat{n}} \mathbf{P}_j = \mathbf{I}$.

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► The spectral projectors give a *resolution of the identity*: $\sum_{j=1}^{n} \mathbf{P}_{j} = \mathbf{I}$.

- $\blacktriangleright \mathbf{P}_{g} := \mathbf{P}_{1} + \cdots + \mathbf{P}_{\widehat{m}}, \qquad \mathcal{U}_{g} := \operatorname{Range}(\mathbf{P}_{g}), \qquad m = \operatorname{dim}(\mathcal{U}_{g}).$
- $\blacktriangleright \mathbf{P}_{b} := \mathbf{I} \mathbf{P}_{g}, \qquad \qquad \mathcal{U}_{b} := \operatorname{Range}(\mathbf{P}_{b}), \qquad \operatorname{dim}(\mathcal{U}_{b}) = n m.$

preliminaries: angles between subspaces

▶ 𝒱 = approximating subspace.

For our problems, $\mathcal{V} = \mathcal{K}_k(\mathbf{A}, \mathbf{x}) := \text{span}\{\mathbf{x}, \mathbf{A}\mathbf{x}, \mathbf{A}^2\mathbf{x}, \dots, \mathbf{A}^{k-1}\mathbf{x}\}.$

 \blacktriangleright \mathfrak{U}_{g} = desired invariant subspace

Measure convergence via the containment gap:

 $\delta(\mathfrak{U}_{g}, \mathfrak{V}) = \max_{u \in \mathfrak{U}_{g}} \sin \angle (\mathbf{u}, \mathfrak{V}) = \max_{u \in \mathfrak{U}_{g}} \min_{\mathbf{v} \in \mathfrak{V}} \frac{\|\mathbf{u} - \mathbf{v}\|}{\|\mathbf{u}\|}.$

• We will monitor how $\delta(\mathfrak{U}_g, \mathfrak{K}_k(\mathbf{A}, \mathbf{x}))$ develops as k increases.

example convergence behavior, $A \neq A^*$



cf. GMRES convergence model of [Nevanlinna 1993] example adapted from [Beattie, E., Rossi 2004]

basic convergence model

Building on [Saad 1980, 1983], [Jia 1995], [Sorensen 2002], [Beattie, E., Rossi 2004], and others....

Theorem [Beattie, E., Sorensen 2005]. Suppose \mathcal{U}_{g} is *reachable* from the Krylov space $\mathcal{K}_{k}(\mathbf{A}, \mathbf{x})$. Then for $k \geq 2m$, $\delta(\mathcal{U}_{g}, \mathcal{K}_{k}(\mathbf{A}, \mathbf{x})) \leq C_{1} C_{2} \min_{\phi \in P_{k-2m}} \max_{z \in \Omega_{b}} |1 - \alpha(z)\phi(z)|.$

- $C_1 = C_1(\mathbf{A}, \mathbf{x})$ = measure of starting vector bias.
- $C_2 = C_2(\mathbf{A}, \Omega_b)$ = measure of eigenvector departure from orthogonality.
- P_{k-2m} = set of polynomials of degree k 2m or less.
- $\blacktriangleright \ \Omega_b \subset \mathbb{C}$ contains the undesired eigenvalues.
- $\alpha(z) = (z \lambda_1) \cdots (z \lambda_m).$

Invariant Subspaces

reachable by

Krylov Subspaces

• If $\mathbf{x} \in \mathbb{C}^n$ lacks a component in the desired eigenvector, e.g.,

$$\mathbf{P}_1 \mathbf{x} = \mathbf{0},$$

the desired eigenvalue/eigenvector is *invisible* to Krylov methods (in exact arithmetic). For example, in the power method

$$\mathbf{A}^{p}\mathbf{x} = \sum_{j=1}^{n} \lambda_{j}^{p} \mathbf{P}_{j}\mathbf{x} = \mathbf{0} + \sum_{j=2}^{n} \lambda_{j}^{p} \mathbf{P}_{j}\mathbf{x},$$

the eigenvalue λ_1 has no influence. (We will address this more later.)

A different problem arises when A has repeated eigenvalues with linearly indpendent eigenvectors (derogatory eigenvalues).

A simple example: $\mathbf{A} = \mathbf{I}$ (identity matrix).

$$\mathfrak{K}_k(\mathbf{A}, \mathbf{x}) = \operatorname{span}\{\mathbf{x}, \mathbf{A}\mathbf{x}, \dots, \mathbf{A}^{k-1}\mathbf{x}\} = \operatorname{span}\{\mathbf{x}\}.$$

The Krylov method converges in one step (happy breakdown), exactly finding one copy of the eigenvalue $\lambda = 1$ and eigenvector **x**.

reachable invariant subspaces

A more perplexing example from Chris Beattie [Beattie, E., Rossi 2004]:

By taking c large, we bias x toward the eigenvector $[0, 1, 0, 0, 0]^T$.

For any *c* the Krylov method breaks down (happily) at iteration k = 3, discovering the Jordan block and invariant subspace

$$\mathbf{V}^* \mathbf{A} \mathbf{V} = \mathbf{S} \begin{bmatrix} 1 & 0 & 0 \\ 1 & 1 & 0 \\ 0 & 1 & 1 \end{bmatrix} \mathbf{S}^{-1}, \qquad \text{Range} \left(\begin{bmatrix} 1 & 0 & 0 \\ c & 1 & 0 \\ 1 & 0 & 0 \\ 1 & 1 & 0 \\ 1 & 1 & 1 \end{bmatrix} \right)$$

The Krylov method finds the 3×3 Jordan block with eigenvector $[0, 0, 0, 0, 1]^T$. Only a set of measure zero x can discover $[0, 1, 0, 0, 0]^T$.

reachable invariant subspaces

• Unlucky choices of $\mathbf{x} \in \mathbb{C}^n$ can (in principle) prevent the Krylov method from seeing a desired (simple) eigenvector.

This behavior is fragile to numerical computations, since *infinitesimal* perturbations to x will add a small component in the desired eigenvector.

Single-vector Krylov methods can (in principle) find one Jordan block associated with each eigenvalue.

This behavior is fragile to numerical computations, since *infinitesimal perturbations split multiple eigenvalues*.

▶ Block Krylov methods (with block size *b*, $\mathbf{X} \in \mathbb{C}^{n \times b}$)

 $\mathfrak{K}_k(\mathbf{A}, \mathbf{X}) = \mathsf{Range}([\mathbf{X} \ \mathbf{A}\mathbf{X} \ \mathbf{A}^2\mathbf{X} \ \cdots \ \mathbf{A}^{k-1}\mathbf{X}]),$

can find *b* linearly independent eigenvectors for a single eigenvalue.

How does the

starting vector x

affect convergence?

Henceforth assume the desired invariant subspace ${\mathfrak U}$ is *reachable* from x:

 $\mathfrak{U}_{g} \subset \mathfrak{K}_{n}(\mathbf{A}, \mathbf{x}).$

How does x influence convergence?

For a single eigenpair $(\lambda_1, \mathbf{u}_1)$ with spectral projector \mathbf{P}_1 , Saad [1980] gives

$$\sin \angle (\mathbf{u}_1, \mathcal{K}_k(\mathbf{A}, \mathbf{x})) \leq \frac{1}{\|\mathbf{P}_1\mathbf{x}\|} \min_{\substack{\phi \in P_{k-1} \\ \phi(\lambda_1) = 1}} \|(\mathbf{I} - \mathbf{P}_1)\psi(\mathbf{A})\|,$$

The leading constant grows as the orientation of x toward u_1 diminishes.

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For *m*-dimensional invariant subspaces \mathcal{U}_g , our bounds replace $1/\|\mathbf{P}_1\mathbf{x}\|$ with

$$C_1 := \max_{\psi \in P_{m-1}} \frac{\|\psi(\mathbf{A})\mathbf{P}_{\mathbf{b}}\mathbf{x}\|}{\|\psi(\mathbf{A})\mathbf{P}_{\mathbf{g}}\mathbf{x}\|} = \max_{\mathbf{v} \in \mathcal{K}_m(\mathbf{A},\mathbf{x})} \frac{\|\mathbf{P}_{\mathbf{b}}\mathbf{v}\|}{\|\mathbf{P}_{\mathbf{g}}\mathbf{v}\|},$$

the ratio of the bad to the good component in the worst approximation to \mathfrak{U}_g from the m-dimensional Krylov space.

 $heta = \angle(\mathbf{x}, \mathfrak{U}_g)$, the angle between \mathbf{x} and its best approximation in \mathfrak{U}_g .



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How do the eigenvalues of A affect convergence?

asymptotic convergence rate determined by eigenvalues

$$\min_{\phi \in P_{k-2m}} \max_{z \in \Omega_{b}} |1 - \alpha(z)\phi(z)|$$

• P_{k-2m} = set of polynomials of degree k - 2m or less.

• $\Omega_b \subset \mathbb{C}$ contains the undesired eigenvalues.

•
$$\alpha(z) = (z - \lambda_1) \cdots (z - \lambda_m)$$

The polynomial approximation problem gives convergence like $C \gamma^k$ for some constant C and rate γ .

- ▶ When $\mathbf{A} = \mathbf{A}^*$, $\Omega_b = [\lambda_{m+1}, \lambda_n]$, and use Chebyshev polynomials to compute the convergence rate γ .
- When Ω_b is a simply connected open subset of \mathbb{C} , use conformal mapping to approach the approximation problem.











Black circles on final figure are Fejér points, asymptotically optimal interpolation points for ϕ .

convergence rate: and granularity of the spectrum

If convergence is very slow, perhaps you are solving the wrong problem.

convergence rate: and granularity of the spectrum

If convergence is very slow, perhaps you are solving the wrong problem. Consider m = 1, where we can use the elementary bound [Saad 1980]

$$\sin \angle (\mathbf{u}_1, \mathcal{K}_k(\mathbf{A}, \mathbf{x})) \leq \frac{1}{\|\mathbf{P}_1\mathbf{x}\|} \min_{\substack{\phi \in P_{k-1} \\ \phi(\lambda_1) = 1}} \|(\mathbf{I} - \mathbf{P}_1)\psi(\mathbf{A})\|.$$

Suppose $\mathbf{A} = \mathbf{A}^*$ and we seek leftmost eigenvalue λ_1 , where

 $\lambda_1 < \lambda_2 \leq \cdots \leq \lambda_n.$

The error bound suggests the progress made at each iteration is like

$$\gamma := rac{\sqrt{\kappa}-1}{\sqrt{\kappa}+1}, \qquad ext{where } \kappa := rac{\lambda_n - \lambda_1}{\lambda_2 - \lambda_1}.$$

When **A** is a discretization of an *unbounded operator*, we expect $\lambda_n = ||\mathbf{A}|| \to \infty$ as $n \to \infty$. The convergence rate goes to one as $n \to \infty$. Thus Krylov subspace methods often perform poorly for PDE eigenvalue problems — *unless the set-up is modified*. Apply the Krylov method to discretizations of the Laplacian in one dimension. How does the convergence rate change as the discretization improves?



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Convergence of Krylov Subspace Projection

The problem becomes immediately apparent if we attempt to run Krylov subspace projection *on the operator itself*,

$$\mathfrak{K}_k(L,f) = \operatorname{span}\{f, Lf, \ldots, L^{k-1}f\}$$

For Lu = -u'' with Dirichlet boundary conditions, u(0) = u(1) = 1, take some starting vector $f \in Dom(L)$, i.e.,

$$f(0) = f(1) = 0.$$

In general $Lf \notin Dom(L)$, so we cannot build the next Krylov direction $L^2f = L(Lf)$. The Krylov algorithm breaks down at the third step.

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The operator setting suggests that we instead apply Krylov to L^{-1} :

$$\mathfrak{K}_k(L^{-1}, f) = \operatorname{span}\{f, L^{-1}f, \dots, L^{-(k-1)}f\}.$$

In this case, L^{-1} is a beautiful compact operator:

$$(L^{-1}f)(x) = \iint f + C_0 + C_1 x_1$$

where we choose C_0 and C_1 so that

$$(L^{-1}f)(0) = (L^{-1}f)(1) = 0.$$

krylov projection applied to the operator

We run the Krylov method on L^{-1} exactly in Mathematica. Denote the eigenvalue estimates at the *k*th iteration as $\theta_1^{(k)} \le \theta_2^{(k)} \le \cdots \le \theta_k^{(k)}$.



Observe superlinear convergence as k increases.

For CG and GMRES applied to operators, see [Winther 1980], [Nevanlinna 1993], [Moret 1997], [Olver 2009], [Kirby 2010]. For "superlinear" convergence in finite dimensional settings, see [van der Sluis, van der Vorst, 1986], [van der Vorst, Vuik, 1992], [Beattie, E., Rossi 2004], [Simoncini, Szyld 2005].

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This mode of computation is preferred for discretization matrices as well: the <u>shift-invert Arnoldi method</u> uses $\mathcal{K}_k((\mathbf{A} - \mu \mathbf{I})^{-1}, \mathbf{x})$.

polynomial preconditioning: a cheap spectral transformation

Replace the conventional Krylov space

$$\mathfrak{K}_k(\mathbf{A}, \mathbf{x}) = \operatorname{span}\{\mathbf{x}, \mathbf{A}\mathbf{x}, \mathbf{A}^2\mathbf{x}, \dots, \mathbf{A}^{k-1}\mathbf{x}\}$$

with the polynomial preconditioned transformation

$$\mathfrak{K}_k(\pi(\mathbf{A}), \mathbf{x}) = \operatorname{span}\{\mathbf{x}, \pi(\mathbf{A})\mathbf{x}, \pi(\mathbf{A})^2\mathbf{x}, \dots, \pi(\mathbf{A})^{k-1}\mathbf{x}\}.$$

[Thornquist 2006], [E., Loe, Morgan arXiv:1806.08020]

Use the polynomial π to separate the interesting eigenvalues.

Often increases matvecs, but decreases iterations (hence orthogonalization).

For example, with Hermitian **A**, take π to be the degree-*d* MINRES residual polynomial [Paige & Saunders 1975], which attains

 $\min_{\substack{p\in P_d\\p(0)=1}} \|p(\mathbf{A})\mathbf{b}\|.$

This polynomial tends to separate smallest-magnitude eigenvalues.

polynomial preconditioning: a cheap spectral transformation

Polynomial preconditioning: Hermitian A, MINRES polynomial.



How do the eigenvectors of A affect convergence?

bounding functions of a matrix

If **A** is normal (**A**^{*}**A** = **AA**^{*}), eigenvectors are orthogonal. For f analytic on σ (**A**), $||f(\mathbf{A})|| = \max_{\lambda \in \sigma(\mathbf{A})} |f(\lambda)|.$

For nonnormal **A**, the situation is considerably more complicated.

► If **A** is diagonalizable,
$$\mathbf{A} = \mathbf{U}\mathbf{A}\mathbf{U}^{-1}$$
, then
 $\|f(\mathbf{A})\| \leq \|\mathbf{U}\| \|\mathbf{U}^{-1}\| \max_{\lambda \in \sigma(\mathbf{A})} |f(\lambda)|.$

For the numerical range (field of values) $W(\mathbf{A}) = \{\mathbf{v}^* \mathbf{A} \mathbf{v} : \|\mathbf{v}\| = 1\},\ \|f(\mathbf{A})\| \leq (1 + \sqrt{2}) \max_{z \in W(\mathbf{A})} |f(z)|.$

► For the *\varepsilon*-pseudospectrum $\sigma_{\varepsilon}(\mathbf{A}) = \{z \in \sigma(\mathbf{A} + \mathbf{E}) \text{ for some } \|\mathbf{E}\| < \varepsilon\},\$ $\|f(\mathbf{A})\| \leq \frac{L_{\varepsilon}}{2\pi\varepsilon} \max_{z \in \sigma_{\varepsilon}(\mathbf{A})} |f(z)|,$

where L_{ε} is the boundary length of $\sigma_{\varepsilon}(\mathbf{A})$.

constant to account for nonnormality

 $C_2 = C_2(\mathbf{A}, \Omega_b)$ comes from bounding $||f(\mathbf{A}|_{\mathcal{U}_b})||$, $f(z) = 1 - \alpha(z)\phi(z)$.

Theorem [Beattie, E., Sorensen 2005]. Suppose \mathcal{U}_{g} is *reachable* from the Krylov space $\mathcal{K}_{k}(\mathbf{A}, \mathbf{x})$. Then for $k \geq 2m$, $\delta(\mathcal{U}_{g}, \mathcal{K}_{k}(\mathbf{A}, \mathbf{x})) \leq C_{1} C_{2} \min_{\substack{\phi \in P_{k-2m} \\ z \in \Omega_{b}}} \max_{z \in \Omega_{b}} |1 - \alpha(z)\phi(z)|.$

▶ If $\Omega_b = \sigma(\mathbf{A}|_{\mathcal{U}_b})$ (no defective eigenvalues), then $C_2 = \|\mathbf{U}_b\| \|\mathbf{U}_b^+\|$, where the columns of $\mathbf{U}_b \in \mathbb{C}^{n \times (n-m)}$ are eigenvectors of $\mathbf{A}|_{\mathcal{U}_b}$.

• If $\Omega_b = W(\mathbf{A}|_{\mathbf{u}_b})$ then $C_2 = 1 + \sqrt{2}$.

• If
$$\Omega_b = \sigma_{\varepsilon}(\mathbf{A}|_{\mathcal{U}_b})$$
 then $C_2 = L_{\varepsilon}/(2\pi\varepsilon)$.

Tension: balance $C_2 \ge 1$ verses size of Ω_b .

transient behavior of the power method

Large coefficients in the expansion of \mathbf{x}_0 in the eigenvector basis can lead to cancellation effects in $\mathbf{x}_k = \mathbf{A}^k \mathbf{x}_0$.

Example: here different choices of α and β affect eigenvalue conditioning,

$$\mathbf{A} = \begin{bmatrix} 1 & \boldsymbol{\alpha} & 0 \\ 0 & 3/4 & \boldsymbol{\beta} \\ 0 & 0 & -3/4 \end{bmatrix}, \quad \mathbf{u}_1 = \begin{bmatrix} 1 \\ 0 \\ 0 \end{bmatrix}, \quad \mathbf{u}_2 = \begin{bmatrix} -4\boldsymbol{\alpha} \\ 1 \\ 0 \end{bmatrix}, \quad \mathbf{u}_3 = \begin{bmatrix} 8\boldsymbol{\alpha}\boldsymbol{\beta}/21 \\ -2\boldsymbol{\beta}/3 \\ 1 \end{bmatrix}.$$



[Trefethen & E. 2005]

restarting

krylov suspace

methods

an essential tool for controlling subspace dimension

restarted Arnoldi algorithm (eigs)

To compute m < k eigenvalues of $\mathbf{A} \in \mathbb{C}^{n \times n}$, Arnoldi methods restrict \mathbf{A} to act on the *k*-dimensional Krylov subspace

$$\mathsf{Ran}(\mathbf{V}) = \mathsf{span}\{\mathbf{x}, \mathbf{A}\mathbf{x}, \dots, \mathbf{A}^{k-1}\mathbf{x}\}.$$

Compute eigenvalues of V^*AV (*Ritz values*), and order them by relevance; e.g., if seeking the rightmost eigenvalue of **A**, let

$$\operatorname{\mathsf{Re}} \theta_1 \geq \operatorname{\mathsf{Re}} \theta_2 \geq \cdots \geq \operatorname{\mathsf{Re}} \theta_k$$

Exact shifts [Sorensen 1992] *restart* the method, attempting to improve **v** with a *polynomial filter* having the "unwanted" Ritz values as roots:

 $\mathbf{x}_{+} = (\mathbf{A} - \theta_{m+1}\mathbf{I})\cdots(\mathbf{A} - \theta_{k}\mathbf{I})\mathbf{x}.$



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Shade $\sim \log_{10}(magnitude of filter polynomial)$

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To understand convergence, one must understand how the Ritz values are distributed over $\sigma(\mathbf{A})$.



Shade $\sim \log_{10}(magnitude of filter polynomial)$

Pushing the language Haim Avron used the previous talk, a standard Krylov method (fixed k) is a sketch-and-solve method, while restarted Krylov methods sketch-to-precondition.

restarted Arnoldi algorithm (eigs)

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- [Sorensen 1992] proved convergence for $A = A^*$.
- The process fails for some A ≠ A* [E. 2009], [Duintjer Tebbens, Meurant 2012].
- Stringent sufficient conditions are known [Carden 2011].

Do nonsymmetric matrices enjoy any kind of interlacing ?

interlacing is a key to convergence theory for $A = A^*$

Cauchy's interlacing theorem assures us that Ritz values cannot bunch up at the ends of the spectrum.



eigenvalues (red lines) and Ritz values (black dots)

interlacing does not hold for $A = A^*$

The absence of interlacing for non-Hermitian problems is the major impediment to a full convergence theory – and is the mechanism that allows the method to fail (in theory).



eigenvalues (red lines) and Ritz values (black dots = real parts)

a pathologically terrible example

A monster, built using the construction of [Duintjer Tebbens, Meurant 2012]:

$$\mathbf{A} = \begin{bmatrix} 1 & 0 & 0 & 0 & 0 & 0 & 0 & -362880 \\ 1 & 2 & 0 & 0 & 0 & 0 & 0 & -1451520 \\ 1 & 3 & 0 & 0 & 0 & 0 & -1693440 \\ & 1 & 4 & 0 & 0 & 0 & -846720 \\ & & 1 & 5 & 0 & 0 & -211680 \\ & & & 1 & 6 & 0 & -28224 \\ & & & 1 & 7 & -2016 \\ & & & & & 1 & -64 \end{bmatrix}, \qquad \mathbf{x} = \begin{bmatrix} 1 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \end{bmatrix}.$$

arXiv:1801.00234



ritz value localization for non-hermitian matrices

Do non-Hermitian matrices obey any kind of "interlacing" theorem?

Ritz values must be contained within the numerical range

 $W(\mathbf{A}) = \{\mathbf{v}^* \mathbf{A} \, \mathbf{v} : \|\mathbf{v}\| = 1\},\$

a closed, convex subset of \mathbb{C} that contains $\sigma(\mathbf{A})$.

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a closed, convex subset of \mathbb{C} that contains $\sigma(\mathbf{A})$.

Consider an extreme example:

$$\mathbf{A} = \begin{bmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 0 & 0 & 0 \end{bmatrix}, \qquad W(\mathbf{A}) = \left\{ z \in \mathbb{C} : |z| \le \frac{\sqrt{2}}{2} \right\}.$$

Repeat the following experiment many times:

• Generate random *two* dimensional subspaces, $\mathcal{V} = \text{Ran } \mathbf{V}$, where $\mathbf{V}^* \mathbf{V} = \mathbf{I}$.

- Form $\mathbf{V}^* \mathbf{A} \mathbf{V} \in \mathbb{C}^{2 \times 2}$ and compute Ritz values $\{\theta_1, \theta_2\} = \sigma(\mathbf{V}^* \mathbf{A} \mathbf{V})$.
- Identify the leftmost and rightmost Ritz values.
- Since $\sigma(\mathbf{A}) = \{0\}$, "interlacing" is meaningless here...

ritz values of a jordan block



ritz values of a jordan block



leftmost Ritz value

rightmost Ritz value

10,000 random (complex) two dimensional subspaces

three matrices with identical W(A)

Compute k = 4 Ritz values for these 8×8 matrices.



 $(\gamma_1 \text{ and } \gamma_3 \text{ set to give same } W(\mathbf{A}) \text{ for all examples; } \varrho = 1/8.)$

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Smallest magnitude of k = 4 Ritz values, 10,000 random complex subspaces.

ritz value localization, sorted by real part

Using Schur's eigenvalue majorization theorem for Hermitian matrices, we can establish an interlacing-type result.

Theorem (Carden & E. 2012)

Let $\theta_1, \ldots, \theta_k$ denote the Ritz values of $\mathbf{A} \in \mathbb{C}^{n \times n}$ drawn from a k < n dimensional subspace, labeled by decreasing real part: $\operatorname{Re} \theta_1 \geq \cdots \geq \operatorname{Re} \theta_k$. Then for $j = 1, \ldots, k$,

$$\frac{\mu_{n-k+j}+\cdots+\mu_n}{k-j+1} \le \operatorname{Re} \theta_j \le \frac{\mu_1+\cdots+\mu_j}{j}$$

where $\mu_1 \geq \cdots \geq \mu_n$ are the eigenvalues of $\frac{1}{2}(\mathbf{A} + \mathbf{A}^*)$.

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where $\mu_1 \geq \cdots \geq \mu_n$ are the eigenvalues of $\frac{1}{2}(\mathbf{A} + \mathbf{A}^*)$.

• The fact that $\theta_i \in W(\mathbf{A})$ gives the well-known bound

$$\mu_1 \leq \operatorname{\mathsf{Re}} heta_j \leq \mu_n, \qquad j=1,\ldots,k.$$

The theorem provides sharper bounds for interior Ritz values.

- The interior eigenvalues of ¹/₂(A + A^{*}) give additional insight; cf. eigenvalue inclusion regions of [Psarrakos & Tsatsomeros, 2012].
- Theorem applies to any subspace Range(V): Krylov, block Krylov, etc.

three matrices with identical W(A)

Three matrices with the same $W(\mathbf{A})$, different interior structure; 2000 trials. For k = 4, numbers on right indicate max Ritz values in each region.



The log-majorization of products of eigenvalues by products of singular values [Marshall, Olkin, Arnold 2011] leads to a limit on Ritz value magnitudes.

Theorem (Carden & E., 2012)

Let $\theta_1, \ldots, \theta_k$ denote the Ritz values of $\mathbf{A} \in \mathbb{C}^{n \times n}$ drawn from a k < n dimensional subspace, labeled by decreasing magnitude: $|\theta_1| \ge \cdots \ge |\theta_k|$. Then for $j = 1, \ldots, k$,

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Related results:

Zvonimir Bujanovic [2011] studies Ritz values of normal matrices *from Krylov subspaces* in his Ph.D. thesis (Zagreb).

Jakeniah Christiansen [2012] studies real Ritz values for n = 3 (SIURO).

some closing thoughts

Krylov methods can further develop as a prominent tool for RandNLA.

- Polynomials are better than powers! Krylov methods have great advantages over power/subspace iteration.
- Block methods hold promise but additional subtleties. Large subspaces can be built rapidly; must maintain linear independence.
- Restarting is crucial in engineering computations, but analysis is tricky. Restarting controls the subspace dimension, refines the starting vector.
- Spectral transformations (shift-invert) can vastly accelerate convergence. You are not entirely constrained by the eigenvalue distribution of A.
- Non-Hermitian problems are solved everyday. The theory is incomplete and monsters are easy to construct, but the Krylov method (as implemented in eigs/ARPACK) works well.
- Can Random Matrix Theory shed light on Ritz value locations? What is the probability that A is stable if V*AV is stable?